# Zero-Temperature Dynamics of Ising Models on the Triangular Lattice 

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Received June 26, 2001; accepted August 10, 2001


#### Abstract

We consider the nature of spin flips of zero-temperature dynamics for ferromagnetic Ising models on the triangular lattice with nearest-neighbor interactions and an initial configuration chosen from a symmetric Bernoulli distribution. We prove that all spins flip infinitely many times for almost every realization of the dynamics and initial configuration.


KEY WORDS: Ising models; triangular lattice; spin flips.

## 1. INTRODUCTION

Let $\mathscr{G}=\left(E_{\mathscr{G}}, V_{\mathscr{G}}\right)$ be an infinite graph with edge set $E_{\mathscr{G}}$ and vertex set $V_{\mathscr{G}}$. Consider the stochastic process $\sigma^{t}=\left\{\sigma_{x}^{t}, x \in V_{\mathscr{G}}\right\}$ on $\mathscr{G}$ which corresponds to the zero-temperature limit of Glauber dynamics for the Ising model with Hamiltonian

$$
H=-\sum_{x, y: x \sim y} J_{x, y} \sigma_{x} \sigma_{y},
$$

where $x \sim y$ denotes that $x$ and $y$ are nearest neighbors, i.e., $\{x, y\} \in E_{\mathscr{G}}$. The process $\sigma^{t}$ takes values in $\Omega=\{-1,+1\}^{V_{s}}$, the space of spin configurations. The initial value $\sigma^{0}=\left\{\sigma_{x}^{0}, x \in V_{g}\right\}$ is taken from a symmetric Bernoulli product measure, i.e., for each $x \in V_{\mathscr{g}}, \sigma_{x}^{0}$ takes +1 (or -1 ) with probability $1 / 2$, independently for different vertex. At each $x$, there is a "Poisson clock" (with rate 1) which "rings" at random times independently

[^0]between different vertices. When the Poisson clock rings at time $t$, the spin at $x$ will consider to flip (i.e., $\sigma_{x}^{t+0}=-\sigma_{x}^{t-0}$ ). If the change in energy
$$
\triangle H_{x}(\sigma)=2 \sum_{y: y \sim x} J_{x, y} \sigma_{x} \sigma_{y}
$$
(where $\sigma=\sigma^{t-0}$ ) is negative (or positive), then the spin at $x$ flips (or does not flip) with probability 1 . If $\triangle H_{x}(\sigma)=0$, then a fair coin is tossed to determine if the spin at $x$ flips. We denote by $P_{\tau}$ the probability distribution on the times at which the Poisson clock rings along with the fair coin tosses, and by $P_{\sigma}{ }^{0} \times \tau=P_{\sigma}{ }^{0} \times P_{\tau}$ the joint distribution of the independent $\sigma^{0}$ and $\tau$.

We are interested in whether for almost every $\sigma^{0}$ and $\tau$ (i.e., $P_{\sigma \times \tau}$-almost surely), $\sigma^{\infty}\left(\sigma^{0}, \tau\right)=\lim _{t \rightarrow \infty} \sigma^{t}\left(\sigma^{0}, \tau\right)$ exists, that is, whether for every $x \in V_{g}$, $\sigma_{x}^{t}\left(\sigma^{0}, \tau\right)$ flips only finitely many times. If $J_{x, y}$ 's are independent random variables and $\mathscr{J}=\left\{J_{x, y},\{x, y\} \in E_{\mathscr{G}}\right\}$ is chosen from the product measure $P_{\mathscr{f}}$ of a continuously distributed probability measure on the real line, then it is proved in ref. 2, under very mild conditions, that for almost every $\mathscr{J}$, $\sigma^{0}$, and $\tau, \sigma^{t}\left(=\sigma^{t}\left(\sigma^{0}, \tau\right)\right)$ flips only finitely many times for every $x \in V_{\mathscr{G}}$, for a very general class of infinite graphs, including the hypercubic lattice $Z^{d}$, the hexagonal and triangular lattice, homogeneous trees, etc. On the other hand, if $J_{x, y}=+1$ (or a positive constant) for all $\{x, y\} \in E_{\mathscr{G}}$, then the picture is less complete. In this case it is proved in ref. 2 that (1) $\sigma_{x}^{t}$ flips infinitely many times for almost every $\sigma^{0}$ and $\tau$ and for every $x \in V_{G}$ if $\mathscr{G}=Z$ or $Z^{2}$, and (2) $\sigma_{x}^{t}$ flips only finitely many times for almost every $\sigma^{0}$ and $\tau$ and for every $x \in V_{\mathscr{G}}$ if $\mathscr{G}$ is a transitive graph in which each vertex has an odd number of nearest neighbors, e.g., the hexagonal lattice and homogeneous trees of odd degrees. See also ref. 1 for the model on the homogeneous tree of degree three. In this note, we consider the model on the triangular lattice and prove the following theorem.

Theorem 1. Let $\mathscr{G}$ be the triangular lattice and $J_{x, y}=+1$ for all $\{x, y\} \in E_{\mathscr{g}}$. Then for almost every $\sigma^{0}$ and $\tau$ and for every $x \in V_{\mathscr{G}}, \sigma_{x}^{t}$ flips infinitely many times.

The idea of the proof of the theorem is from that of Theorem 2 in ref. 2. The key in the proof is the translation ergodicity along the three lines passing through each vertex. It is unknown if $\sigma_{x}^{t}$ flips infinitely many times or only finitely many times when $\mathscr{G}$ is $Z^{d}$ with $d \geqslant 3$, or a homogeneous tree of even degrees, or a hyperbolic graph (see ref. 3 for a definition) in which each vertex has an even number of nearest neighbors.

## 2. PROOF OF THE THEOREM

First of all, since the distributions $P_{\tau}$ and $P_{\sigma^{0}}$ are translation-invariant, so is the product distribution $P_{\tau, \sigma^{0}}=P_{\tau} \times P_{\sigma}{ }^{0}$. For a given vertex $x$, let $A_{x}^{+}$ (or $A_{x}^{-}$) be the event that $\sigma_{x}^{\infty}\left(\tau, \sigma^{0}\right)$ exists and equals +1 (or -1 ), and denote by $I_{x}^{+}$(or $I_{x}^{-}$) the indicate function of this event. By translationinvariance and symmetry under the global spin flip $\sigma^{0} \rightarrow-\sigma^{0}$, it follows that for all $x, P_{\tau, \sigma^{0}}\left(A_{x}^{+}\right)=P_{\tau, \sigma^{0}}\left(A_{x}^{-}\right)=p$ for some $p \in[0,1 / 2]$. We wish to prove that $p=0$.

Suppose $p>0$. Fix a vertex $a$ in the lattice. By translation-ergodicity, for each of the three lines passing through $a$ there are, with $P_{\tau, \sigma^{\circ}}$-probability one, infinitely many vertices $x$ on each side of $a$ such that $A_{x}^{-}$occurs. This implies that there exists vertices $b, d$ and $f$ on the three lines passing through $a$ as shown in Fig. 1 such that

$$
P_{\tau, \sigma^{0}}\left(A_{a}^{+} A_{b}^{-} A_{d}^{-} A_{f}^{-}\right)>0 .
$$

So there exists some $t_{0}$ such that, with strictly positive $P_{\tau, \sigma^{0}}$-probability, $\sigma_{a}^{t}=+1$ and $\sigma_{b}^{t}=\sigma_{d}^{t}=\sigma_{f}^{t}=-1$ for all $t \geqslant t_{0}$. But this would at least require that the transition probabilities of the Markov process $\sigma^{t}$ satisfies

$$
\inf _{\sigma \in \Omega^{\prime}} P_{\tau}\left(\sigma^{t+1} \notin \Omega^{\prime} \mid \sigma^{t}=\sigma\right)=0,
$$

where $\Omega^{\prime}$ is the set of spin configurations on the triangular lattice such that $\sigma_{a}=+1$ and $\sigma_{b}=\sigma_{d}=\sigma_{f}=-1$. Next, we will reach a contradiction by showing

$$
\begin{equation*}
\inf _{\sigma \in \Omega^{\prime}} P_{\tau}\left(\sigma^{t+1} \notin \Omega^{\prime} \mid \sigma^{t}=\sigma\right)>0 . \tag{2.1}
\end{equation*}
$$

This will prove that $p=0$.
Let $R$ be the finite region enclosed in the polygon "abcdef." More precisely, $R$ is the set of vertices inside (and on the boundary of) the polygon abcdef. Let $\Omega_{R}=\{-1,+1\}^{R}$ be the set of spin configurations in $R$ and $\Omega_{R}^{\prime}$ be the subset of $\Omega_{R}$ with $\sigma_{a}=+1$ and $\sigma_{b}=\sigma_{d}=\sigma_{f}=-1$. Because there are only finitely many elements in $\Omega_{R}^{\prime}$, in order to show (2.1) it is sufficient to show that for each $\sigma_{R}^{\prime} \in \Omega_{R}^{\prime}$

$$
\begin{equation*}
P_{\tau}\left(\sigma^{t+1} \notin \Omega^{\prime}\left|\sigma^{t}\right|_{R}=\sigma_{R}^{\prime}\right)>0, \tag{2.2}
\end{equation*}
$$

where $\left.\sigma\right|_{R}$ is $\sigma$ restricted on $R$.
For any $\sigma_{R}^{\prime} \in \Omega_{R}^{\prime}$, we call a Peierls contour (in the dual lattice) which separates plus and minus spins in $R$ a domain wall. We claim that for any


Fig. 1.
$\sigma_{R}^{\prime} \in \Omega_{R}^{\prime}$, there exists a domain wall which either (1) starts from line segment $a f$ and ends on line segment $a b$, or (2) starts from line segment $a b$ and ends on line segment $b c$, or (3) starts from line segment $a d$ and ends on line segment $c d$, or (4) starts from line segment $a d$ and ends on line segment $d e$, or else (5) starts from line segment $a f$ and ends on line segment ef. To prove the claim, first notice that the spins at vertices $a$ and $f$ are respectively +1 and -1 . So along the line segment from $a$ to $f$, there exists the first vertex at which the spin is -1 . Call this vertex $i_{1}$ and the vertex right before it $i_{1}-1$. Then there is a domain wall starting at the midpoint between vertices $i_{1}-1$ and $i_{1}$. If this domain wall comes back to line segment $a f$ and ends at the midpoint between some $i_{2}-1$ and $i_{2}$, then the spin at $i_{2}$ must be +1 . Call the first vertex after $i_{2}$ which has spin -1 as $i_{3}$. Then there is a domain wall starting at the midpoint between $i_{3}-1$ and $i_{3}$. If this domain wall ends on line segment $a f$ again, then repeat the procedure. Eventually, we can find the first (counting from vertex $a$ to vertex $f$ ) domain wall which starts on line segment $a f$ and ends on either line segment $a b$ or $b c$ or $c d$ or $d e$ or $e f$. If it ends on line segment $a b$ or $e f$, then the proof is finished. If it ends on line segment $b c$, then it cuts region $R$ into two parts: $R_{b}$, the part which contains vertex $b$, and $R-R_{b}$. Notice that the spins on the boundary of $R_{b}$ along this domain wall are all +1 , so any other domain wall in $R_{b}$ can not intersect this domain wall. Find the first domain wall (counting from vertex $b$ to vertex $a$ ) which starts from line segment $b a$ and does not end on $b a$. Then this domain wall must end on $b c$ because it can not end on $a f$, since if it ends on $a f$, then the domain wall which starts from $a f$ and ends on $b c$ would not be the first one not ending on $a f$ (counting from vertex $a$ to vertex $f$ ). Therefore, there exists a domain wall which starts on line segment $a b$ and ends on line segment $b c$. If the domain wall starting on line segment $a f$ ends on line segment $c d$,
then it crosses $a d$ and hence there is a domain wall which starts on $a d$ and ends on $c d$. Finally if the domain wall starting on $a f$ ends on $d e$, then it cuts the parallelogram "adef" into two parts: $R_{d}$, the part which contains vertex $d$, and the rest of the diamond. Because of the way the domain wall is chosen, in $R_{d}$ the vertex $d$ is surrounded by +1 spins, ranging from vertex $a$ to the beginning of the domain wall and then along the domain wall to the end of the domain wall on line segment $d e$. So there must be a domain wall starting on line segment $a d$ and ending on line segment $d e$. This completes the proof of the claim.

We now turn to the proof of (2.2). Without loss of generality, suppose there exists a domain wall which starts on line segment $a f$ and ends on line segment $a b$ (see Fig. 1). The other four cases listed in the claim can be argued similarly. This domain wall cuts $R$ into two parts: $R_{a}$, the part which contains vertex $a$, and $R-R_{a}$. We first give an order to the vertices in $R_{a}$ according to the following rule: order the vertex which is farthest from line $a f$ first; if two or more vertices have the same distance from line $a f$, then order the one which is farthest from line $a b$ first. Then with positive $P_{\tau}$-probability there is some sequence of clock rings (according to the order prescribed above) and coin toss outcomes within $R$ that will move the domain wall towards vertex $a$ so that $\sigma_{a}^{t+1}=-1$. For example, in Fig. 1, the clocks rings in the order of $1,2,3,4,5,6,9$. After the clock at 4 rings, the domain wall encloses only vertices 5,6 , and 9 . From here it is further moved towards vertex $a$ and finally the spin at $a$ is changed from +1 to -1 . This completes the proof of the theorem.

## ACKNOWLEDGMENTS

This work was initiated when the author was visiting the Courant Institute of Mathematical Sciences. He thanks Chuck Newman and the Institute for the warm hospitality he received. He is grateful to Douglas Howard and Chuck Newman for introducing to him this topic. Research was supported in part by NSF Grant DMS-98-03598.

## REFERENCES

1. C. D. Howard, Zero-temperature Ising spin dynamics on the homogeneous tree of degree three, J. Applied Probability 37:736-747 (2000).
2. S. Nanda, C. M. Newman, and D. L. Stein, Dynamics of Ising spin systems at zero temperature, in On Dobrushin's Way (From Probability Theory to Statistical Physics), Amer. Math. Soc. Transl. Ser. 2, Vol. 198 (Amer. Math. Soc., Providence, Rhode Island, 2000), pp. 183-194.
3. R. Rietman, B. Nienhuis, and J. Oitmaa, The Ising model on hyperlattices, J. Phys. A. Math. Gen. 25:6577-6592 (1992).

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